UNIQUENESS OF SOLUTION FOR A CLASS OF STEFAN PROBLEMS

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Abstract

This paper deals with a theoretical mathematical analysis of one-dimensional solidification problem, in which kinetic undercooling is incorporated into the This temperature condition at the interface. A model problem with nonlinear kinetic law is considered. We prove a local result intimate for the uniqueness of solution of the corresponding free boundary problem.

Keywords: Free boundary problem; Nonlinear integral equation

Introduction

It is well known that in many industrial areas, the solidification process plays a significant role. Mathematical models of solidification including interface kinetics effects have been considered for quite some time (see [1], and references therein). This class of free boundary problems, which arises in a number of physical situations, is that of on equilibrium problems, in which the phase - change temperature is dependent on the velocity of the front at which the phase-change occurs (for more physical problems, see [3-7]). Here, we study a model problem with nonlinear kinetic law at the interface in the one-dimensional case. Specifically, let the curve with s(0)=b(0 < b < 1) be defined as the interface that separates the liquid and solid phases. With u denoting temperature (scaled so that is vanishes at equilibrium), we may write the system of equations as

$$u_t = K_l u_{xx}$$
 in $Q_l = \{(x,t) | 0 < x < s(t), 0 < t \le T\},$ (1.1)

$$u_t = K_s u_{xx}$$
 in $Q_1 = \{(x,t) | 0 < x < s(t), 0 < t \le T\},$ (1.2)

and on the interface x = s(t) as

$$u_1 = u_2 = g_1(V(f)), \tag{1.3}$$

$$Ku_{x}^{+} - Ku_{x}^{-} = g_{2}(V(t)), \qquad (1.4)$$

$$s(0) = b, \ 0 < b < 1,$$
 (1.5)

where K_l and K_s are thermal diffusivities of a liquid and a solid respectively, L > 0 is the latent heat and the superscripts + and – denote, respectively the right-hand and left-hand limits with respect to the special variable x. These equations are subject to the initial and boundary conditions

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$$u(x,0) = \varphi_1(x) \quad 0 \le x \le b$$
, (1.6)

$$u(x,0) = \varphi_2(x) \quad b \le x \le 1$$
, (1.7)

$$u(i-1,t) = f_i(t) \quad t \ge 0, \quad (i=1,2)$$
 (1.8)

where

$$V(t) = \frac{ds(t)}{dt}$$
(1.9)

is the propagation velocity of the free boundary. The free boundary problem considered here was formulated in [1], where reduction the problem to an integral equation was given. In the context of solid fuel combustion, s(t) represents the boundary between the unburnt and burnt material, and u_1, u_2 , are the nondimensionalized temperature in the unburnt and burnt material respectively, (see [3-7] and references therein). The temperature at the free boundary controls its velocity $V(t) = g_1^{-1}(u_1(s(t), t))$. The heat exchange between the unburnt (x < s(t)) and burnt material is modeled by the boundary condition in (1.4) which, in principle, may be nonlinear.

Main Results

Theorem. Consider the problem (1.1)-(1.9). Suppose that the kinetic function and initial and boundary data satisfy the assumption $(H_1) - (H_3)$ in [1]. Then the problem (1.1)-(1.9) has not more than one solution.

To prove uniqueness for $t < \sigma$ suppose that $u_0 = (u_{01}, u_{02}), s_0$ is another solution of (1.1)-(1.9) for $t < \sigma$ and $v_0(t) = (v_{01}(t), v_{02}(t))^T$ is another solution of integral equations (26) and (27) in [1]. It suffices to prove uniqueness, for any $\overline{\sigma} < \sigma$.

Let

$$\overline{M} = Max\{M, l.u.b_{0 \le t \le \overline{\sigma}} |v_0(t)|\}$$

where M introduced in section 4.2 in [1], and let be any positive number satisfying

$$\sqrt{\sigma} < Min \left\{ \left(C_2 \| \varphi_1^{\cdot} \| + C_3 \| f_1^{\cdot} \| + C_4 + C_5 \overline{M} + C_6 + C_7 \right)^{-1} \frac{\overline{M}}{2}, \\ \left(D_2 + D_3 + D_4 \overline{M} + D_5 + D_6 \| \varphi_2^{\cdot} \| + D_7 \| f_2^{\cdot} \| \right)^{-1} \frac{\overline{M}}{2} \right\}$$

where the constants C_i and D_i , i=2,3,...,7 are simple combination of $\pi, b, \frac{1}{h}, M, M', M_2, M'_2, K$. Then by the same calculations in [1] which were used to prove that T maps $B_{M,\sigma}$ into itself (where T and $B_{M,\sigma}$ introduced in subsection 4.2 in [1]) and is a contraction one shows that T maps $B_{\overline{M},\overline{\sigma}}$ into itself and is a contraction. Hence, there exists at most one fixed point of T in $B_{\overline{M},\overline{\sigma}}$. It follows that $v(t) = v_0(t)$ for $0 \le t \le \overline{\sigma}$, where v(t) is solution of integral equations (26) and (27) in [1]. Hence also $s(t) = s_0(t)$ $u(x,t) = u_0(x,t)$ if $0 \le t \le \overline{\sigma}$, $0 \le x \le s(t)$ and $s(t) \le x \le 1$. We next consider the system (1.1)-(1.9) for $t > \sigma$, i.e. (1.1)-(1.5), (1.8), (1.9) are considered for $t > \overline{\sigma}$ (instead of $t \ge 0$) where as (1.6), (1.7) are replaced by $u_1(x,\overline{\sigma}) = u_1(x,\overline{\sigma})$ for $0 \le x \le s(\overline{\sigma})$, $u_2(x,\overline{\sigma}) = u_2(x,\overline{\sigma})$ for $s(\overline{\sigma}) < x < 1$.

This problem can again be transformed into integral equations (26), (27) in [1] extend to the present integral equation provided M is replaced by M_0 where

$$M_0 = \underbrace{lub}_{\sigma < t < \sigma} \quad \left| V(t) g_1(V(t)) \right|$$

Similarly to section 4 in [1], we reduce the problem (1.1)-(1.9) for u_0, s_0 in the interval $\overline{\sigma} \le t < \sigma$ to an integral equation. Since $u_1(x, \overline{\sigma}) = u_1(x, \overline{\sigma})$, $u_2(x, \overline{\sigma}) = u_2(x, \overline{\sigma})$, the integral equation for v(t) and $v_0(t)$ coincide. Repeating now the same argument as before we conclude that for $v(t) = v_0(t)$ for any $\overline{\sigma}$ satisfying

$$\begin{split} &\sqrt{\left(\widetilde{\sigma}-\overline{\sigma}\right)} < \\ &Min\left\{\left(C_{2}\left\|\varphi_{1}^{'}\right\|+C_{3}\left\|f_{1}^{'}\right\|+C_{4}+C_{5}\overline{M_{0}}+C_{6}+C_{7}\right)^{-1}\frac{\overline{M_{0}}}{2}, \\ &\left(D_{2}+D_{3}+D_{4}\overline{M_{0}}+D_{5}+D_{6}\left\|\varphi_{2}^{'}\right\|+D_{7}\left\|f_{2}^{'}\right\|\right)^{-1}\frac{\overline{M_{0}}}{2}\right\} \end{split}$$

We can now proceed in the same manner as before in [2] step by step, nothing that in each step the time interval can be taken to be $\geq \varepsilon$ where satisfies

$$\sqrt{\varepsilon} < Min\left\{ \left(C_2 \| \varphi_1^{\cdot} \| + C_3 \| f_1^{\cdot} \| + C_4 + C_5 \overline{M_1} + C_6 + C_7 \right)^{-1} \frac{\overline{M_1}}{2} \right\}$$

$$\left(D_2 + D_3 + D_4 \overline{M_1} + D_5 + D_6 \| \varphi_2^{\cdot} \| + D_7 \| f_2^{\cdot} \| \right)^{-1} \frac{\overline{M_1}}{2} \right\}$$

where

$$\overline{M}_{1} = Max\{ \underset{\overline{\sigma} < t < \sigma}{lub} |V(t)g_{1}(V(t))|, \underset{\overline{\sigma} < t < \sigma}{lub} |v_{0}(t)| \}$$

Having proved existence and uniqueness for all $t < \sigma$ where σ is any positive number satisfying (36) in [1]. Let us stress that the previous proof (see (38), (39) in [1]) shows also the following:

If instead of (1.1)-(1.9) for t > 0 we consider (1.1)-(1.9) for $t > \lambda$, i.e., (1.1)-(1.5), (1.8), (1.9) hold for $t > \lambda$ and (1.6), (1.7) replaced by $u_1(x, \lambda) = u_1(x, \lambda)$ for $0 < x < s(\lambda)$ and $u_2(x, \lambda) = u_2(x, \lambda)$ for $s(\lambda) < x \le 1$ respectively, and if

$$|V(\lambda)g_1(V(\lambda))|, \ s(\lambda), \ \frac{1}{s(\lambda)}$$

are bounded independently of λ , then there exists a unique solution for the problem in an interval $\lambda \le t \le \lambda + \varepsilon$, where ε is some positive number independent of λ .

Since for any solution of (1.1)-(1.9) the function s(t) is monotone non-decreasing, $\frac{1}{s(\lambda)} \leq \frac{1}{b}$. To complete the proof of theorem it suffices to prove the following statement:

For every $t_0 > 0$ there exists an $\varepsilon > 0$ such that if the system (1.1)-(1.9) has a unique solution for all $t < t_0$, then it also has a unique solution for all $t < t_0 + \varepsilon$ in view of the previous remarks it suffices to show: If

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u(x,t), s(t) is a solution of (1.1)-(1.9) for all $t < t_0$, then for all $\eta > 0$ sufficiently small, the functions

$$|l.u.b|V(t_0 - \eta)g_1(V(t_0 - \eta))|, \ s(t_0 - \eta)$$
(2.1)

are bounded independently of η . If we prove that

 $l.u.b|v(t)| < \infty$,

then from (28) in [1] follows the boundedness of s(t) for $t < t_0$. Consequently, if we prove (2.2) then the proof of theorem is completed.

Proof of (2.2). We use for v(t) the integral equation which corresponds to the system (1.1)-(1.9) in the interval $t_0 - \mu < t < \mu$ (μ sufficiently small) in [1]. Since

$$u_1(0,t_0-\mu) = f_1(t_0-\mu), u_2(0,t_0-\mu) = f_2(t_0-\mu),$$

the equations are

 v_1

$$\begin{split} &(t) = -g_1(V(t))V(t) + \\ & 2 \int_0^{s(t_0 - \mu)} u_{1\xi}(\xi, t_0 - \mu)N(s(t), t; \xi, t_0 - \mu)d\xi \\ & + \int_{t_0 - \mu}^t v_1(\tau)G_x(s(t), t; s(\tau), \tau)d\tau \\ & - 2 \int_{t_0}^t f_1'(\tau)N(s(t), t; 0, \tau)d\tau + \\ & 2g_1(V(t_0 - \mu))N(s(t), t; s(t_0 - \mu), t_0 - \mu) \\ & + 2 \int_{t_0 - \mu}^t g_1'(V(\tau))N(s(t), t; s(\tau), \tau)d\tau \\ & - 2 \int_{t_0 - \mu}^t g_1(V(\tau))V(\tau)G_x(s(t) - 1, t; s(\tau), \tau)d\tau \end{split}$$

$$\begin{split} v_{2}(t) &= -g_{1}(V(t))V(t) \\ &- 2g_{1}(V(t_{0} - \mu))N(s(t) - 1, t; s(\tau) - 1, t_{0} - \mu) \\ &+ \int_{t_{0} - \mu}^{t} g_{1}'(V(\tau))N(s(t) - 1, t; s(\tau) - 1, \tau)d\tau \\ &- 2\int_{t_{0} - \mu}^{t} v_{2}(t)G_{x}(s(t) - 1, t; s(\tau) - 1, \tau)d\tau \\ &+ 2\int_{s(t_{0} - \mu)}^{1} u_{2\xi}(\xi, t_{0} - \mu)N(s(t) - 1, t; \xi, t_{0} - \mu)d\xi \\ &+ \int_{t_{0}}^{t} f_{2}'(\tau)N(s(t) - 1, t; 0, \tau)d\tau \end{split}$$

In section 4 in [1] we proved $v_1(t)$ and $v_2(t)$ are bounded functions, we obtain that

$$\lim_{0 < t < t_0} \left| v(t) \right| < \infty$$

therefore we established theorem.

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